

On a problem in the Stability Discussion of Rotating black holes

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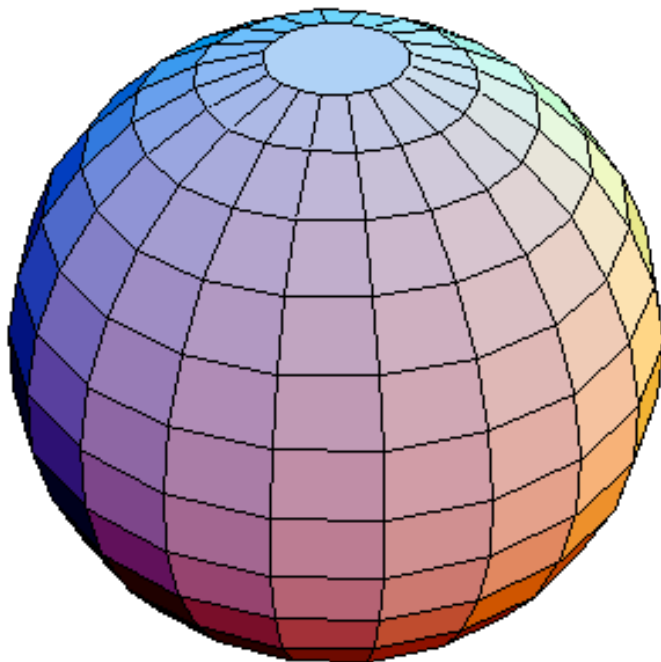
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What is a metric?

- Roughly speaking, a metric is a function ds^2 describing the distance between two points in a set of events S



- Distance = the length of the shortest curve between two points
- The metric of the unit sphere

$$X(\theta, \phi) = (\sin\theta \cos\phi, \sin\theta \sin\phi, \cos\theta)$$

(arc length on a meridian):

$$d\Omega^2 = d\theta^2 + \sin^2\theta d\phi^2$$



Einstein's field equations

- 10 nonlinear PDEs
- Foundation of general relativity
- Describe the gravitational effects produced by a mass on the space around it (how it “curves” spacetime)
- Solutions = “metrics” of spacetime
- Minkowski metric
- Schwarzschild metric
- Kerr metric



cdt Kerr black holes

- A more complex solution to EFE, discovered by Roy Kerr in 1963, the Kerr metric describes the geometry of spacetime around a rotating massive body
- Kerr black holes (rotating black holes) believed to be the most frequent in nature, since most stars that undergo gravitational collapse are rotating
- The ultimate question posed by the Kerr metric:
is it stable under gravitational perturbations?
- Any body orbiting a Kerr black hole, distorting its gravitational field, is an example of a perturbation



cdt Kerr black holes

- The stability of the Kerr metric assumed to be true; not mathematically proven
 - Physical observations of black hole behavior
 - A black hole formed in a cluster of stars eventually starves
 - No observational or numerical record of a black hole behaving otherwise



The Klein-Gordon equation on a Kerr background

- Stability under gravitational perturbations → evolution of a gravitational field propagating in the gravitational field of a Kerr black hole
- Gravitational fields - spin 2
- Scalar fields - simplest fields, spin 0
- Current studies focused on the case of the scalar field
- The (reduced) Klein-Gordon equation on a Kerr background
- Describes a spin-0 quantum field (pion field) propagating in the gravitational field of a rotating black hole
- Stability of the Kerr metric \Leftrightarrow non-exponentially growing solutions

$$(\square + \mu^2)u = 0$$



The Klein-Gordon equation on a Kerr background

- Pions \leftarrow the Klein-Gordon field \leftarrow dictated by the black hole's gravitational field \leftrightarrow the coefficients of the \square - operator



The \square - operator

$$\square = \frac{\gamma}{\Delta\Sigma} \partial_t^2 + \frac{4imMar}{\Delta\Sigma} \partial_t - \frac{1}{\Sigma} \partial_r \Delta \partial_r - \frac{1}{\Sigma \sin \theta} \partial_\theta \sin \theta \partial_\theta - \frac{m^2}{\Sigma} \left(\frac{a^2}{\Delta} - \frac{1}{\sin^2 \theta} \right)$$

$$\Delta = r^2 + a^2 - 2Mr$$

$$\Sigma = r^2 + a^2 \cos^2 \theta$$

$$\gamma = \Sigma(r^2 + a^2) + 2Mra^2 \sin^2 \theta$$

- t - time
- r - radial coordinate
- θ - angular coordinate
- μ - mass of solution field
- m - angular momentum # of the solution
- M - mass of black hole
- i - imaginary unit
- a - rotational parameter



Multiplication of Operators

$$\frac{\partial}{\partial x} \left(xy^2 \frac{\partial}{\partial y} \right) = y^2 \frac{\partial}{\partial y} + xy^2 \frac{\partial^2}{\partial x \partial y}$$

$$O_1 = xy \frac{\partial}{\partial x}$$

$$O_2 = xy^2 \frac{\partial}{\partial y}$$

$$O_1 O_2 = xy^3 \frac{\partial}{\partial y} + x^2 y^3 \frac{\partial^2}{\partial x \partial y}$$

$$O_2 O_1 = xy^3 \frac{\partial}{\partial x} + x^2 y^3 \frac{\partial^2}{\partial x \partial y}$$



Carter and Lenaghan's symmetry operator

$$\bar{\square} = \frac{1}{\Sigma} \left\{ a^2 \left[\cos^2 \theta \frac{(r^2 + a^2)^2}{\Delta} + r^2 \sin^2 \theta \right] \partial_t^2 + 2ima \left(r^2 + a^2 \cos^2 \theta \frac{r^2 + a^2}{\Delta} \right) \partial_t - m^2 \left(\frac{a^4 \cos^2 \theta}{\Delta} + \frac{r^2}{\sin^2 \theta} \right) - a^2 \cos^2 \theta \partial_r \Delta \partial_r + \frac{r^2}{\sin \theta} \partial_\theta \sin \theta \partial_\theta \right\}$$

- Carter and Lenaghan's result (1979)
- Commutes with \square
- Presence of 2nd order time derivatives in the \square operator poses a structural problem; natural ideas which emerge:
 - Find a symmetry (commuting?) operator which only contains first order time derivatives
 - Use \square to find this new operator \rightarrow replace the 2nd-rder time derivative in the \square operator with the expression for this derivative obtained from the Klein-Gordon equation



CCT Results

$$\hat{\square} = \frac{2ima(r^2 + a^2)\Sigma}{\gamma} \partial_t + \frac{a^2 \Delta \sin^2 \theta}{\gamma} \partial_r \Delta \partial_r + \frac{(r^2 + a^2)^2}{\gamma \sin \theta} \partial_\theta \sin \theta \partial_\theta + \frac{m^2 [a^4 \sin^4 \theta - (r^2 + a^2)]}{\gamma \sin^2 \theta}$$

$$\hat{\square} \square - \square \hat{\square} = A \square + B \partial_r \square + C \partial_\theta \square$$



Results

$$C = \frac{8Mr a^2 (r^2 + a^2)^3 \sin \theta \cos \theta}{\Sigma \gamma^2}$$

$$B = \frac{-4M \Delta a^2 \sin^2 \theta [2r^2 a^2 \Delta \sin^2 \theta + \Sigma(r^4 - a^4)]}{\Sigma \gamma^2}$$

$$A = \frac{1}{\gamma^3 \Sigma} \left\{ 2Ma^2 \sin^2 \theta (T_1 T_2 - \gamma \Delta T_3) + 4Mr a^2 (r^2 + a^2)^3 [(2 \cos^2 \theta - \sin^2 \theta) \gamma + 4 \Delta a^2 \sin^2 \theta \cos^2 \theta] \right\}$$

$$T_1 = 2r^2 a^2 \Delta \sin^2 \theta + \Sigma(r^4 - a^4)$$

$$T_2 = (2r - 2M) \gamma + 4r \Delta (r^2 + a^2) - 4M(r^4 - a^4)$$

$$T_3 = 4ra^2 \Delta \sin^2 \theta + 2r^2 a^2 (2r - 2M) \sin^2 \theta + 2r(r^4 - a^4) + 4r^3 \Sigma$$



Additional Results

- Initial data \rightarrow numerically \rightarrow solution to the Klein-Gordon equation
- $\bar{\square}$ \rightarrow constraint \rightarrow conserved quantity
- Currently no interpretation (angular momentum?)

$$\square' = a^2 \sin^2 \theta \partial_t^2 + 2i m a \partial_t + \frac{1}{\sin \theta} \partial_\theta \sin \theta \partial_\theta - \frac{m^2}{\sin^2 \theta}$$

$$\square' \square - \square \square' = A'(\square) + B'(\partial_\theta \square)$$

$$A' = \frac{2a^2(2\cos^2 \theta - \sin^2 \theta)}{\Sigma}$$

$$B' = \frac{4a^2 \sin \theta \cos \theta}{\Sigma}$$



ct Broader Impact

- Numerical simulations of collisions between two black holes
 - assume the Kerr metric to be stable
 - show no instability
 - if it were to be proven that the Kerr metric is unstable - why do these simulations show no sign of instability?
 - Numerical dissipation in a finite difference approximation
 - Relatively short evolution times of typical numerical relativity codes
 - Recently improved numerical relativity codes are able to evolve multiple orbits of black hole binaries
 - Equations considered in this theoretical analysis are the EFE *linearized* about the Kerr black hole; numerical simulations solve the full *nonlinear* equations



Broader Impact - LIGO

- LIGO (Laser Interferometer Gravitational-Wave Observatory)
 - Detection of cosmic gravitational waves and their scientific study
 - Strongest sources of such waves - black hole collisions
 - Towards the end phase of the collision - solutions approximated by small perturbations of the Kerr metric (extraction of wave signals from the end phase of a black hole collision)
 - Studies of stellar populations → estimates of the expected number, frequency and location of black hole collision events in space
 - Einstein's theory → limits on the amplitude of gravitational waves detected by LIGO
 - No unexpected events in detected data



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